

Numerical Computation of Black Hole Spacetimes

Mark A. Scheel
Caltech

Outline:

- Introduction
- How are Einstein's equations solved?
- Why is it so difficult to evolve black holes?

Goal

Solve Einstein's equations for binary black hole system

- Dynamical time scale for a single BH is its mass M .
(Reminder: $G = c = 1$. $M_{\odot} = 5\mu s = 1.5\text{km}$)
- Must follow evolution for thousands of M .

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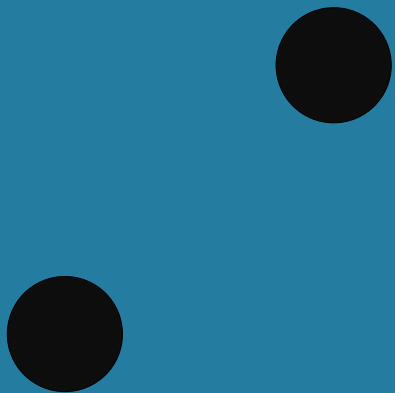
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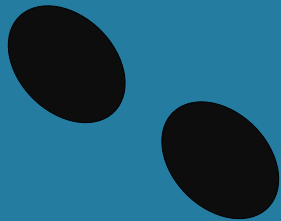
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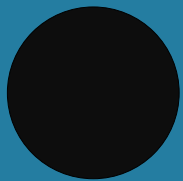


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 - ★ Single Kerr BH
 - ★ Single distorted BH
 - ★ Single boosted BH
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 - ★ “Cosmic screw”
 - ★ Grazing collision
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 - ★ Close quasiequilibrium orbits
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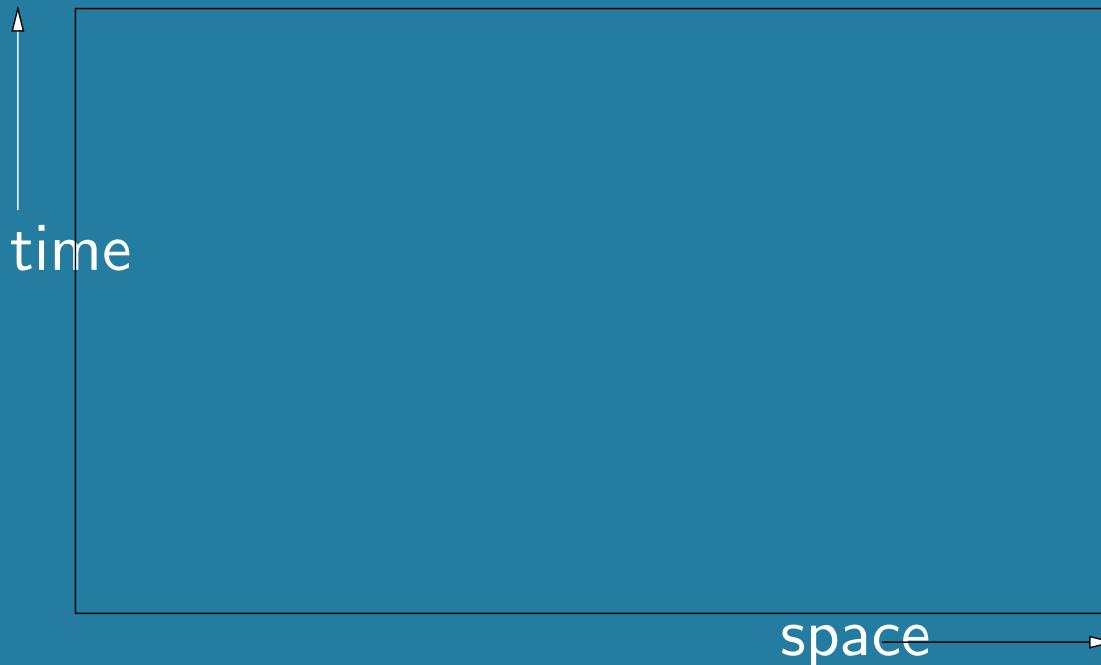
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- All current general methods are unstable!
 - ★ Codes crash at $t \sim 30M$ — $1000M$

3+1 Decomposition

- Key quantity: spacetime metric $g_{\mu\nu}$

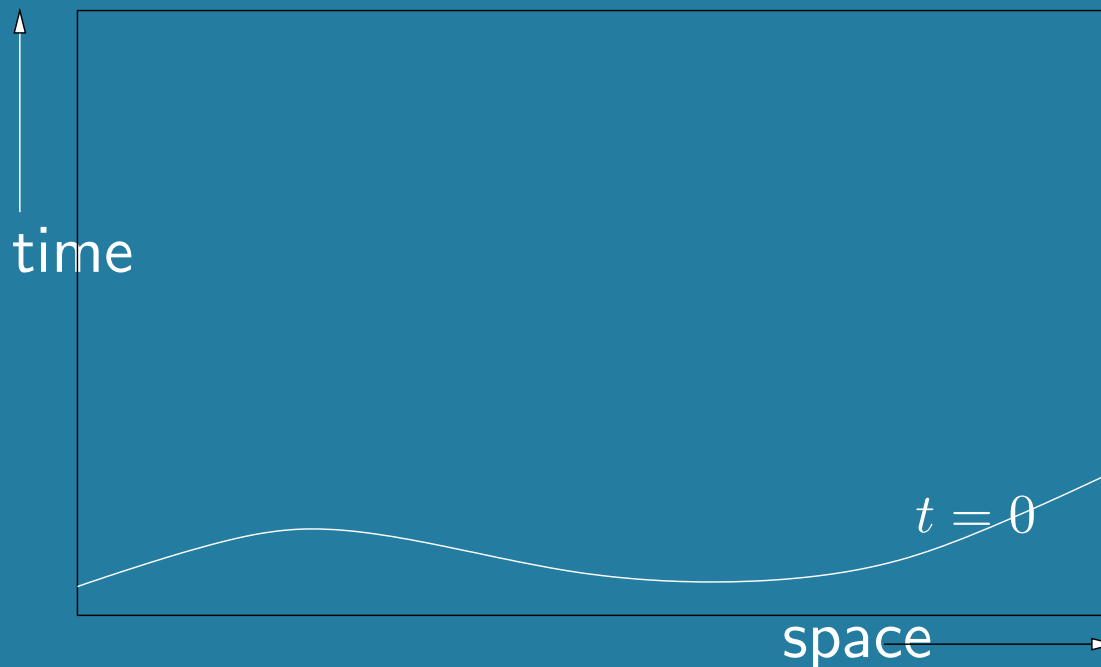
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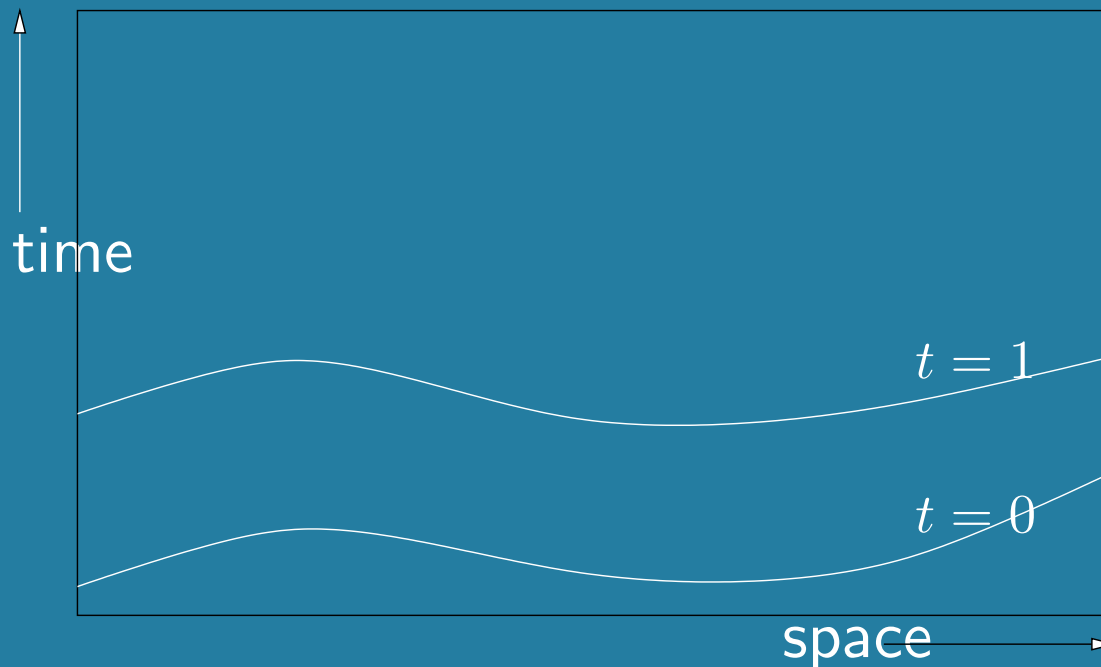
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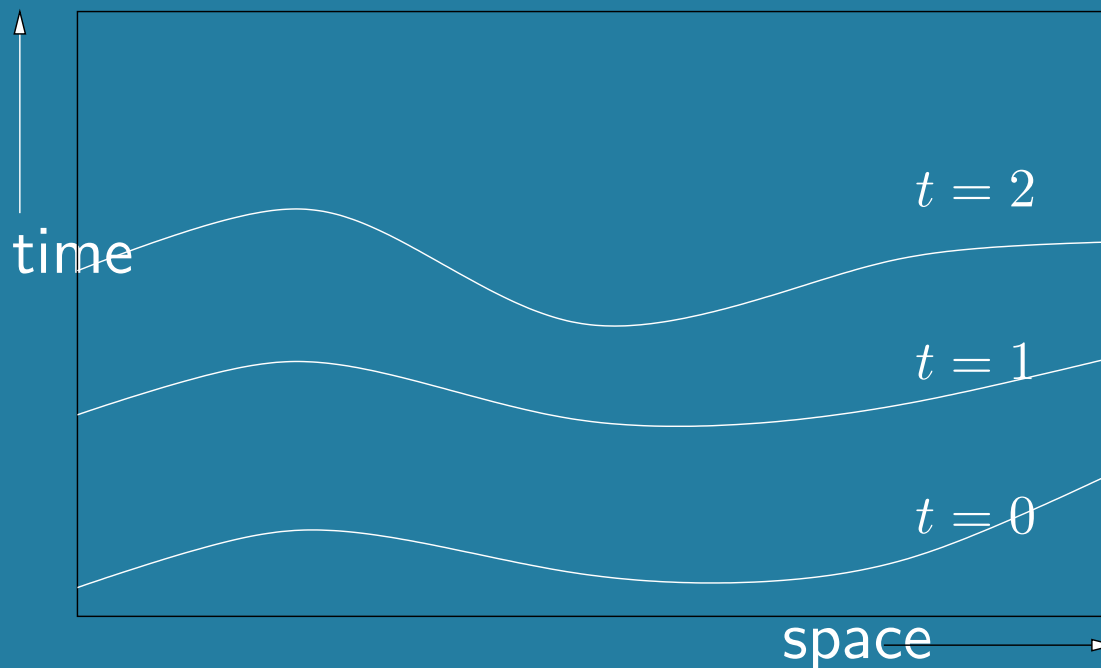
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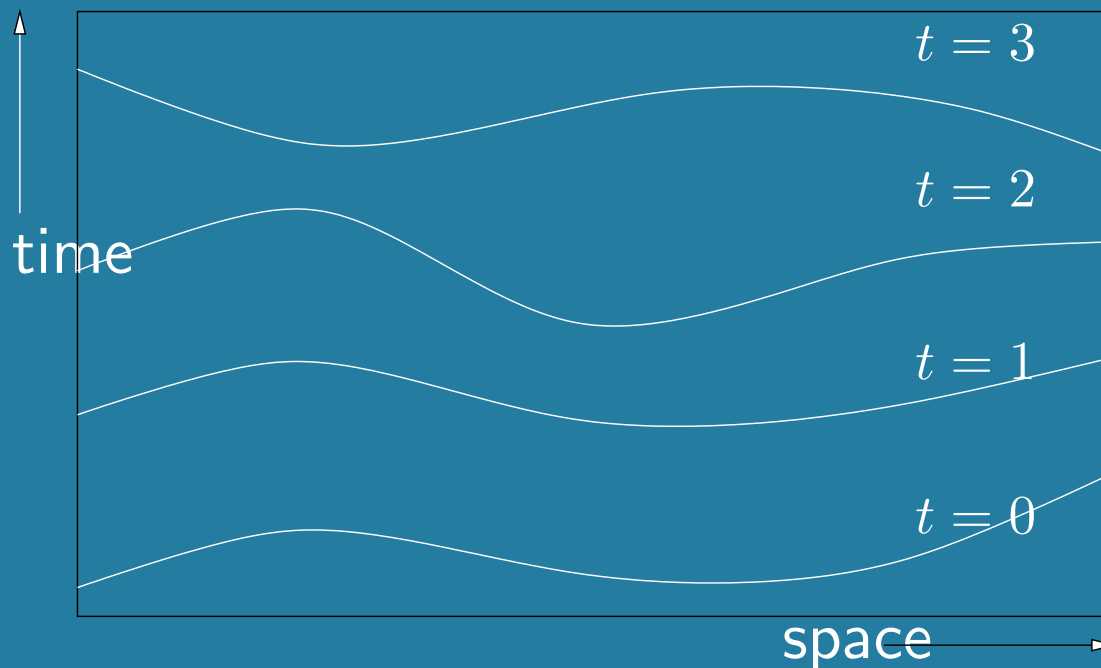
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3+1 Electromagnetism

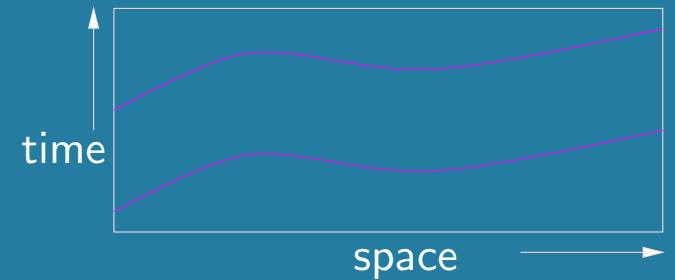
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$$\begin{aligned} A_\mu &\equiv (-\phi, \vec{A}) \\ \vec{E} &\equiv -\partial_t \vec{A} - \vec{\nabla} \phi \end{aligned}$$

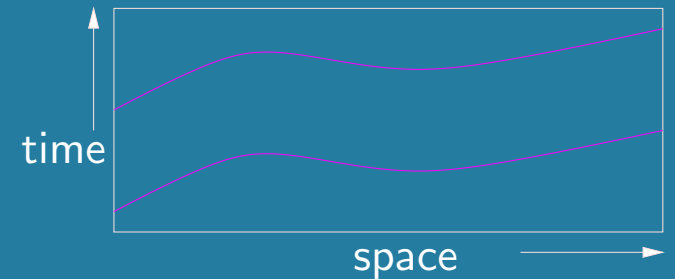


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$$\vec{\nabla} \cdot \vec{E} = 4\pi \rho \quad (C1)$$

• Then $\partial_t \vec{A} = -\vec{E} - \vec{\nabla} \phi \quad (E1)$

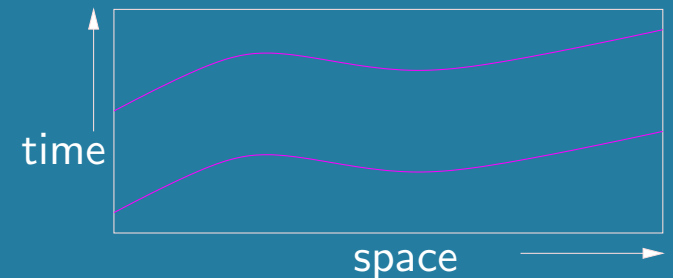
$$\partial_t \vec{E} = \vec{\nabla} \times (\vec{\nabla} \times \vec{A}) - 4\pi \vec{J} \quad (E2)$$

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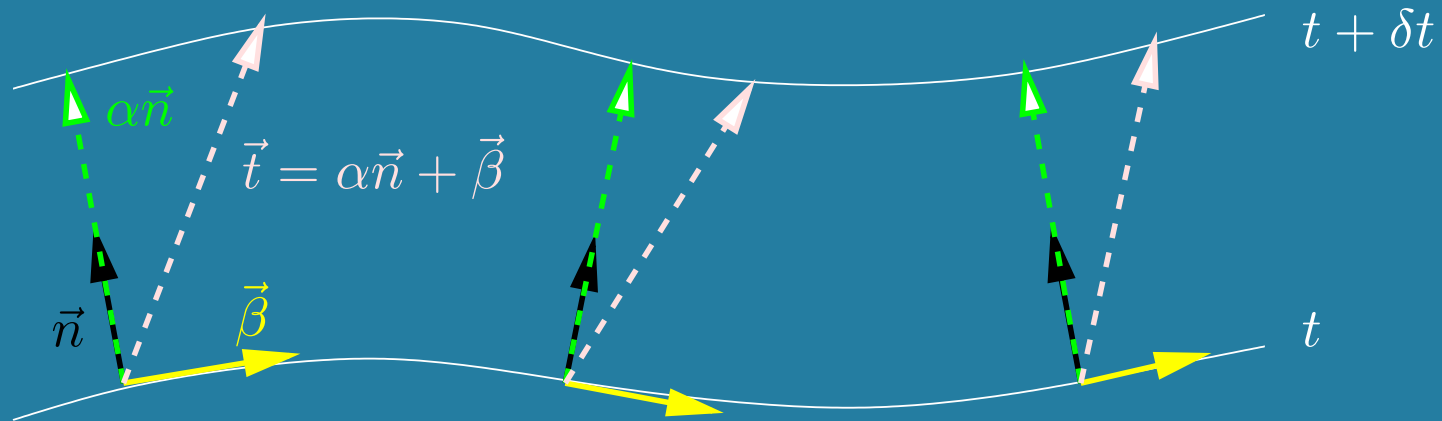
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- Scalar potential ϕ undetermined (gauge).
- Vector potential \vec{A} still has some gauge freedom.

3+1 Variables for General Relativity

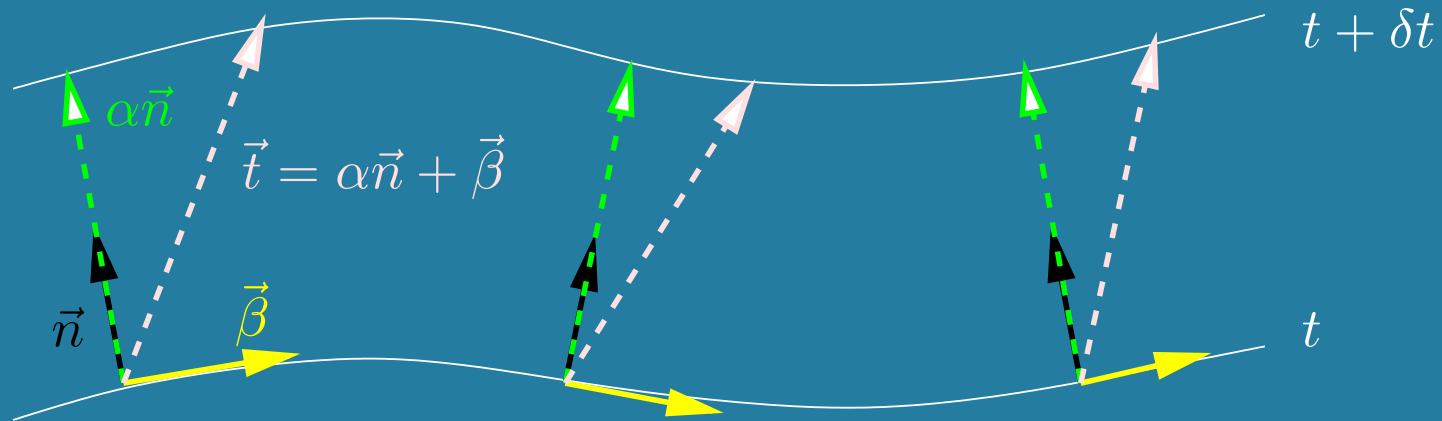
$$ds^2 = -\alpha^2 dt^2 + g_{ij}(dx^i + \beta^i dt)(dx^j + \beta^j dt)$$



- Three-metric g_{ij} : Measures lengths at constant t . (Analogue: \vec{A})
- Lapse α : Measures proper time of normal observer. (Analogue: ϕ)
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- Extrinsic Curvature $K_{ij} \equiv -\frac{1}{2}\mathcal{L}_{\vec{n}}g_{ij}$
Measures how normals diverge. (Analogue: \vec{E})

3+1 Einstein Equations (ADM)

Constraint Equations:

$$\begin{aligned} D^j K_{ij} - D_i K^j_j &= 8\pi S_i \\ \bar{R}^i_i + K^i_i K^i_i - K_{ij} K^{ij} &= 16\pi\rho \end{aligned}$$

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Evolution Equations:

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Maxwell–Einstein Analogy

	Maxwell	Einstein
Variables	A_μ	$g_{\mu\nu}$
Equation	$\partial^\nu F_{\mu\nu} = 4\pi J_\mu$	$G_{\mu\nu} = 8\pi T_{\mu\nu}$
3+1	$A_\mu \rightarrow \vec{A}, \phi$	$g_{\mu\nu} \rightarrow g_{ij}, \alpha, \beta^i$
∂_t variables	\vec{E}	K_{ij}
Constraints	$\vec{\nabla} \cdot \vec{E} = 4\pi\rho$	$F(K, g, K', g', g'') = 0$
Evolution	$\partial_t \vec{A}, \partial_t \vec{E}$	$\partial_t g_{ij}, \partial_t K_{ij}$
Gauge ¹	$A_0 \equiv \phi$	$g_{0\mu} \equiv \alpha, \beta^i$
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¹ Completely undetermined by equations of motion.

² Contains some gauge freedom.

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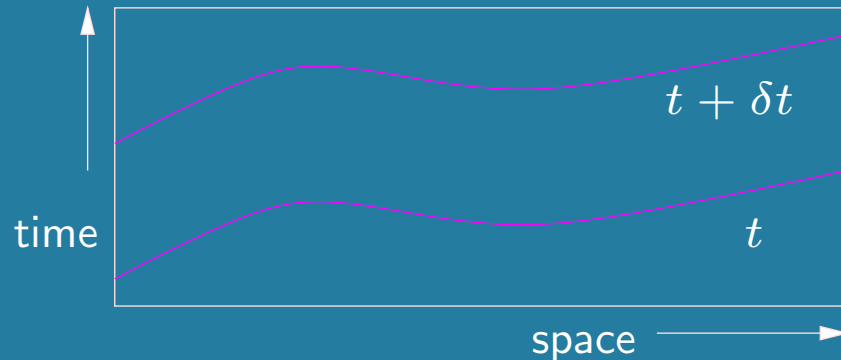
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3+1 Equations Overdetermined

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Maxwell Illustration: $\partial_t \vec{A} = -\vec{E} - \vec{\nabla}\phi \quad (E1)$

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- Given \vec{E} , \vec{A} at t , how does one get \vec{E} , \vec{A} at $t + \delta t$?
- Usually choose **free evolution**.

Why so difficult?

- Complicated equations
- Complicated geometry
- Large CPU/memory requirements

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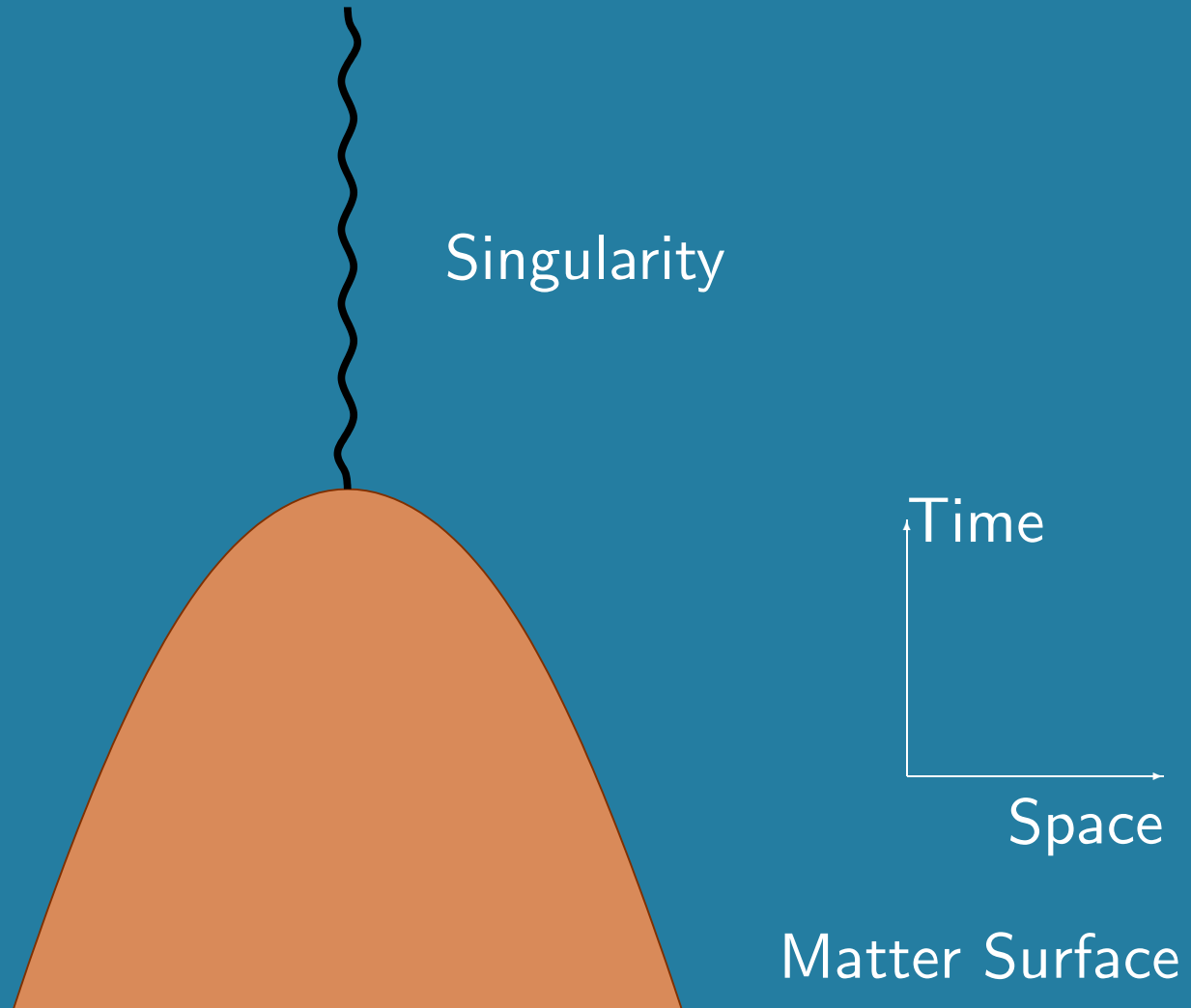
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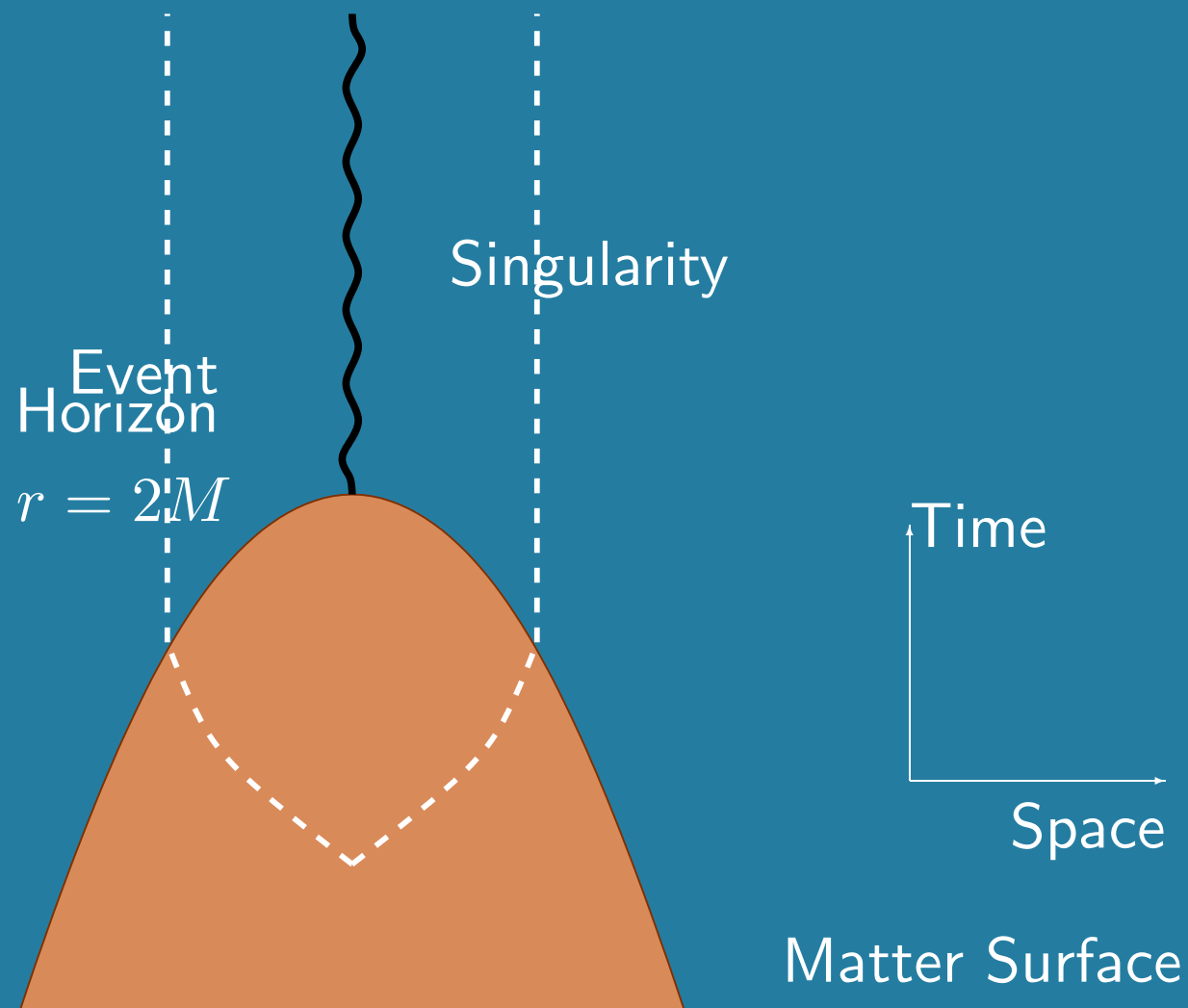
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Bigger computers will **not** solve our problems!

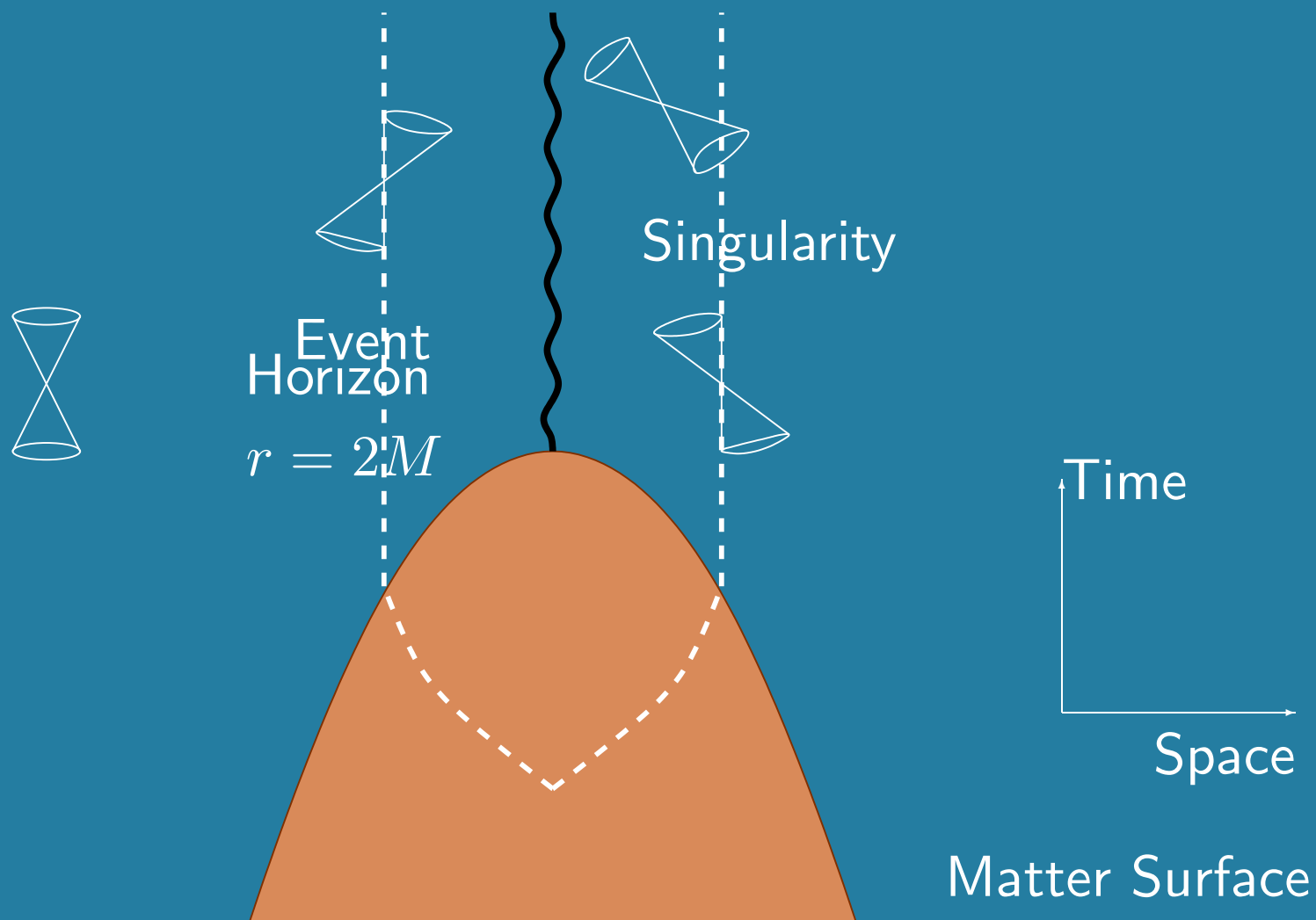
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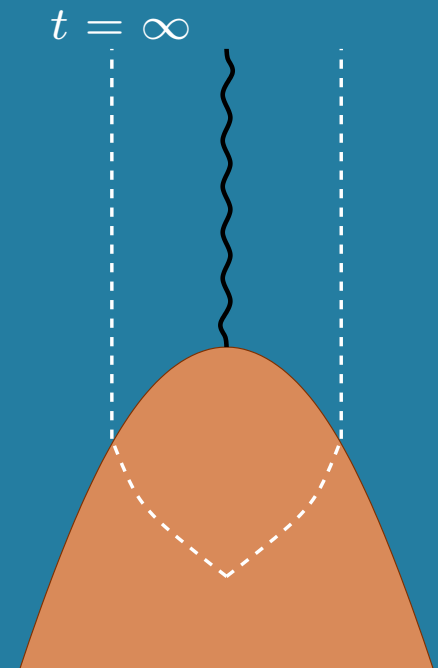


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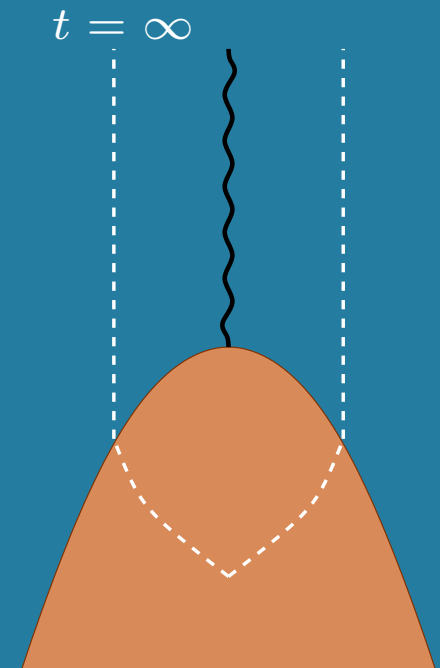
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 - ★ Many bad ones (e.g. Schwarzschild)

$$ds^2 = -(1 - 2M/r) dt^2 + (1 - 2M/r)^{-1} dr^2 + r^2 d\Omega^2$$



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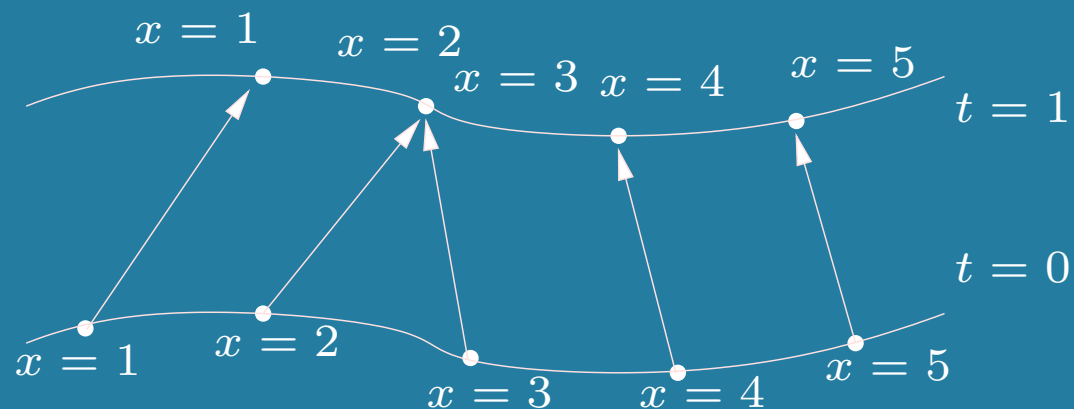
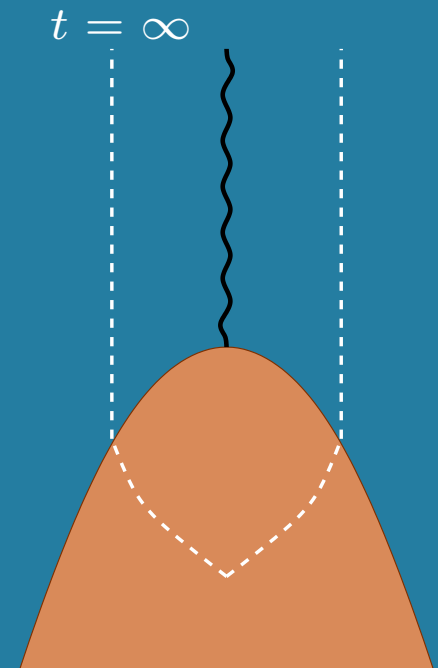
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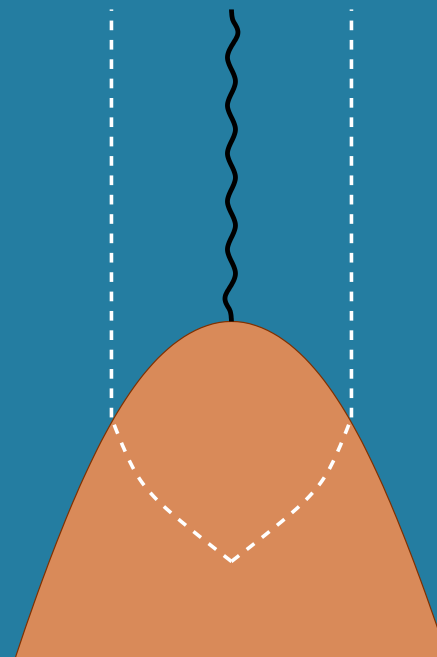
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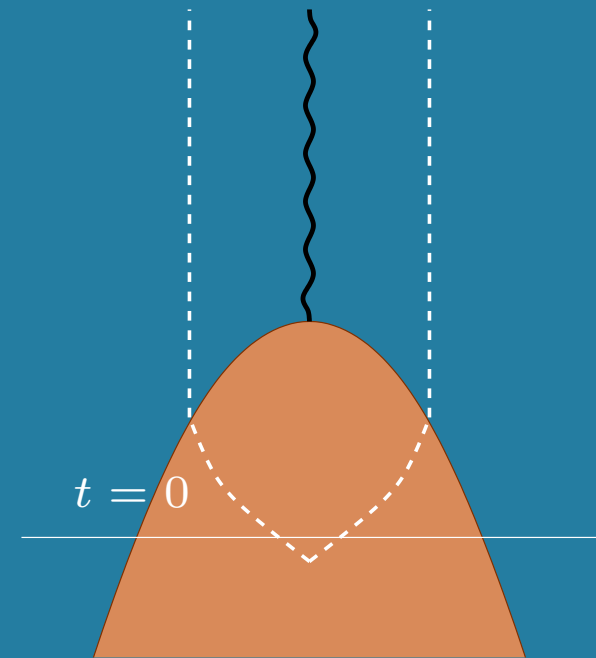


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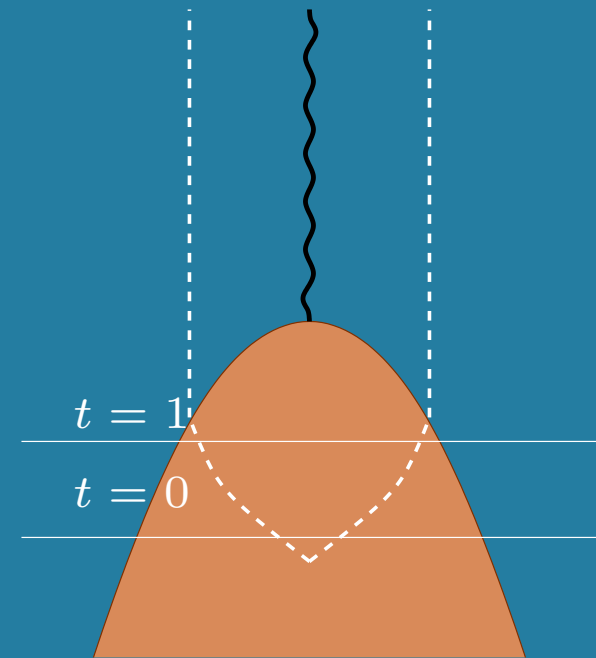
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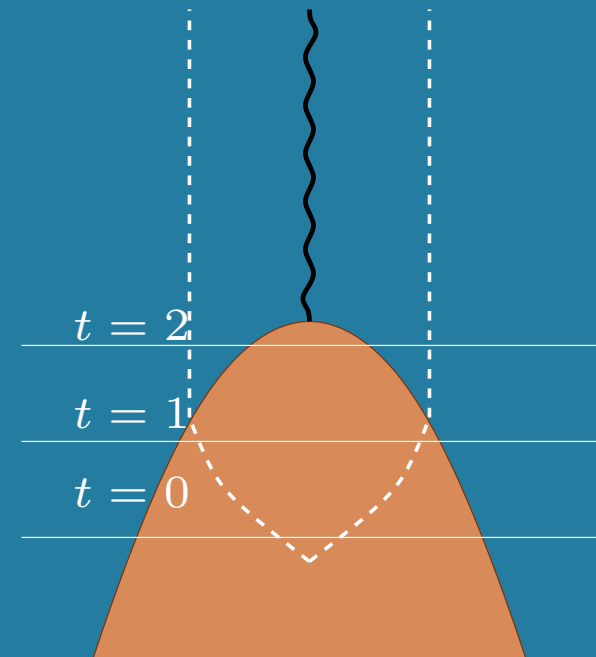
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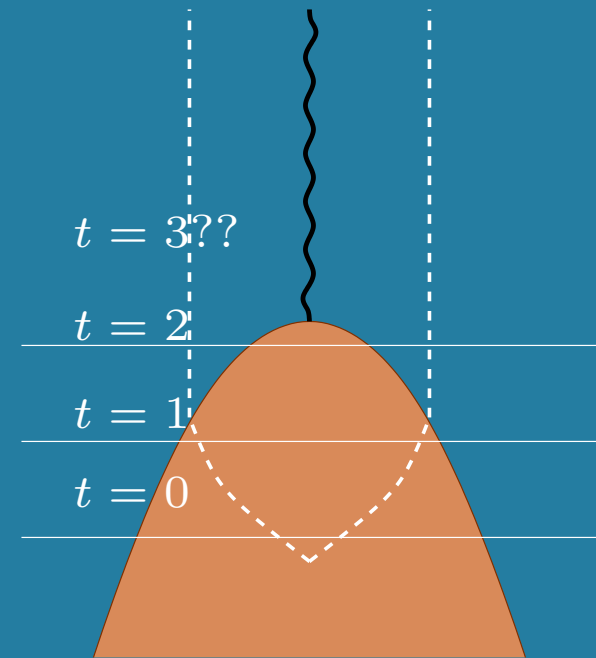
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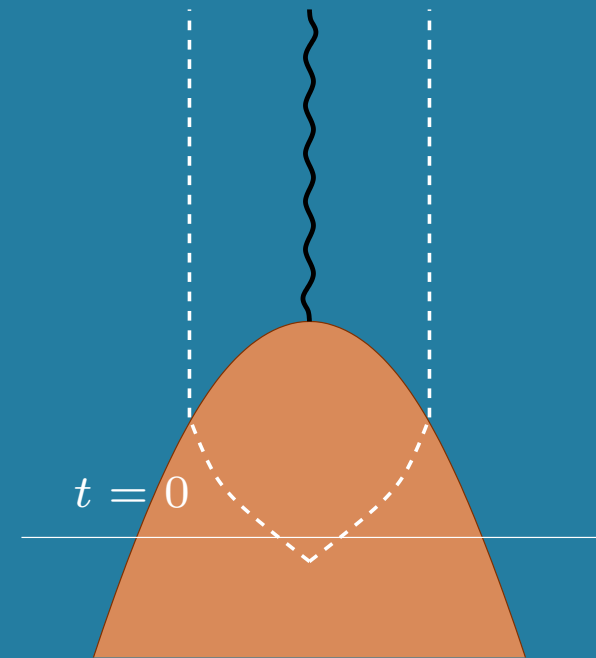
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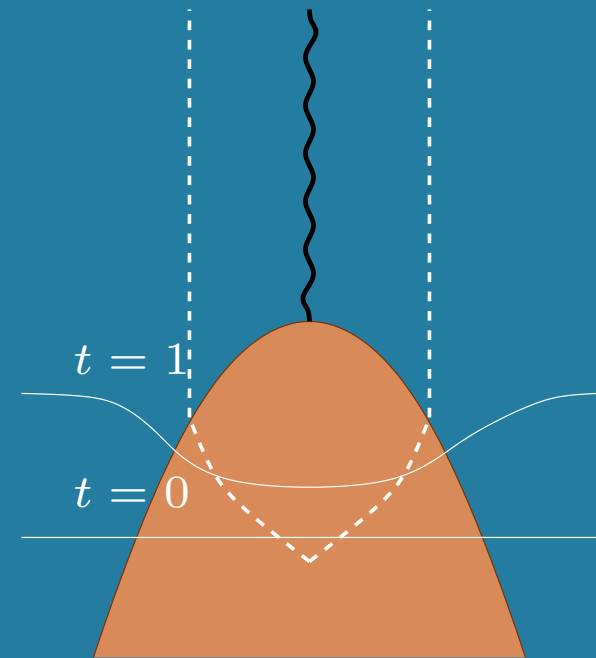
Singularities

- Naive integration hits singularity
- Try coordinate freedom



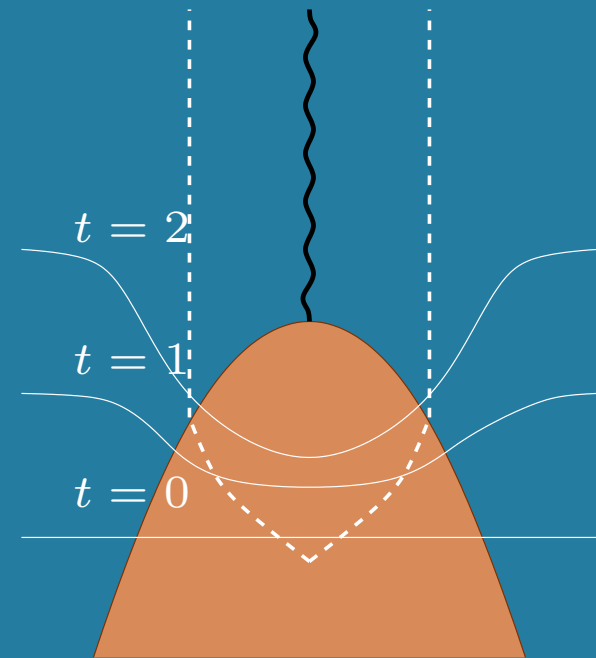
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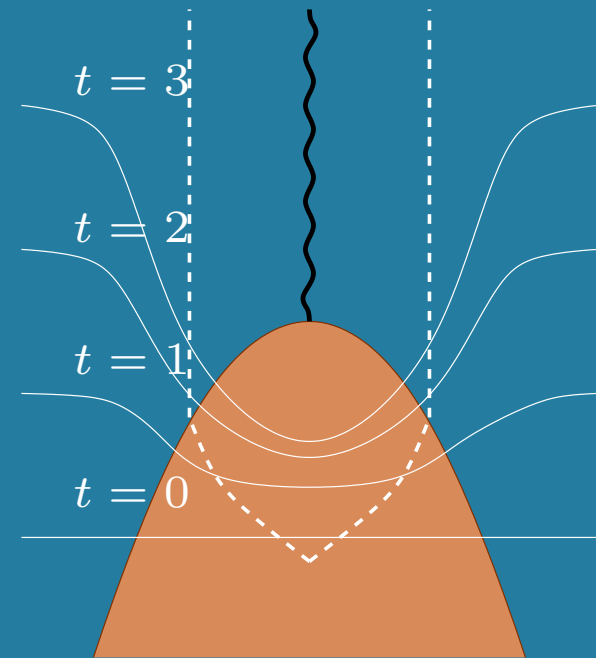
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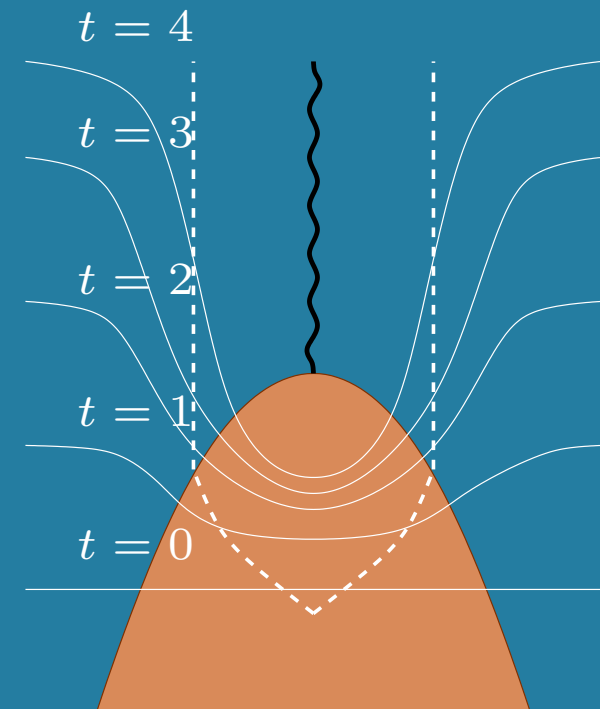
Singularities

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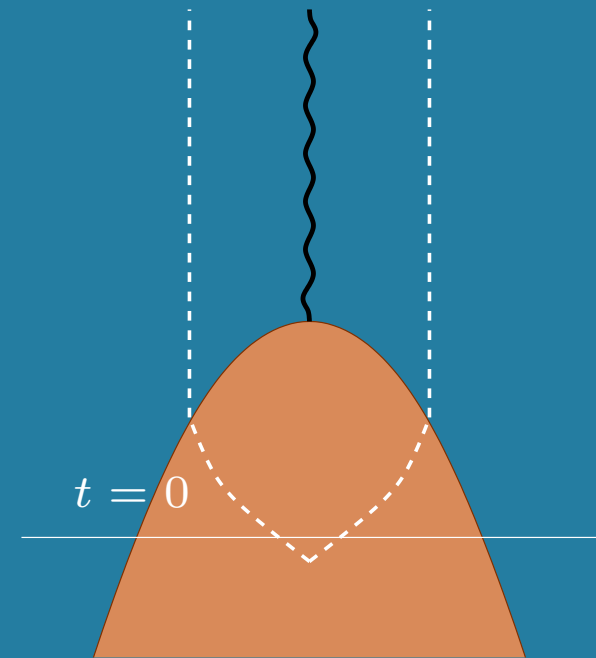
Singularities

- Naive integration hits singularity
- Try coordinate freedom
 - ★ Grid stretching



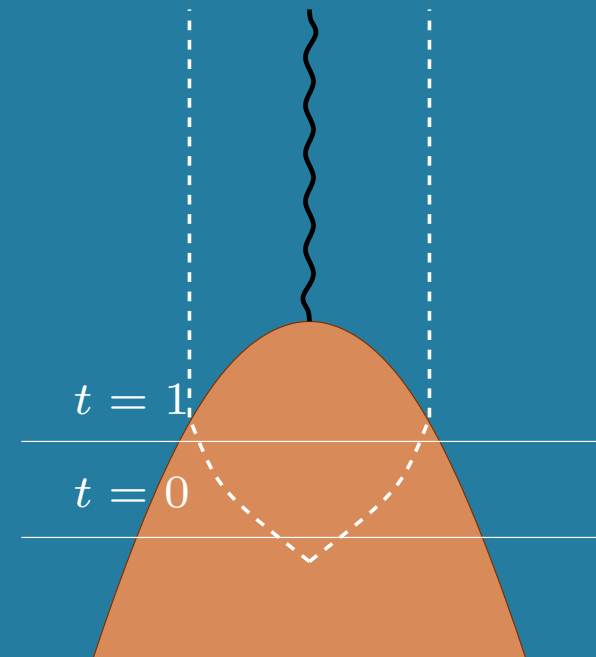
Singularities

- Naive integration hits singularity
- Try coordinate freedom
 - ★ Grid stretching
- Modern method: **Excision**



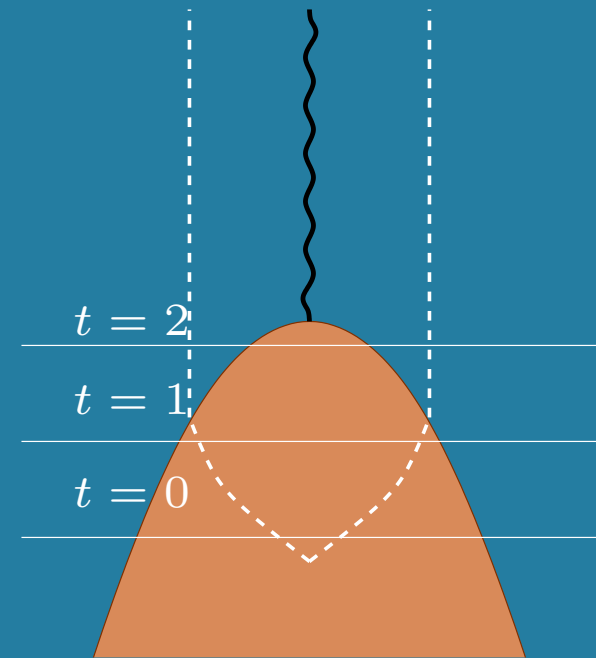
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- Try coordinate freedom
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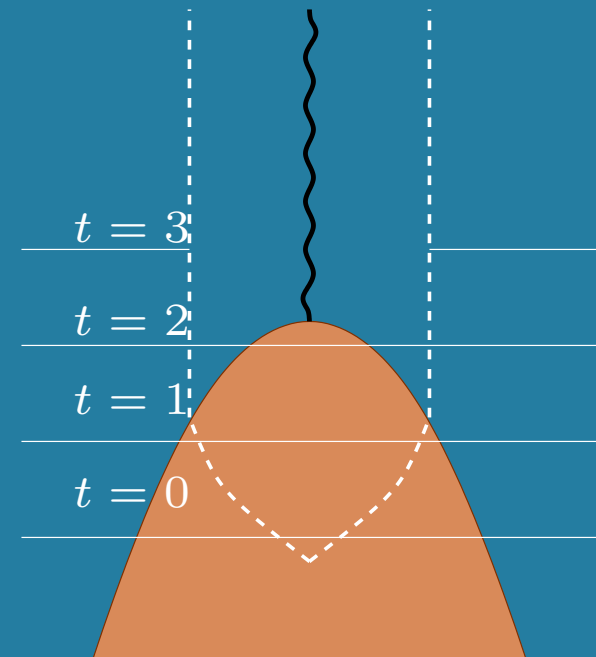
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- Try coordinate freedom
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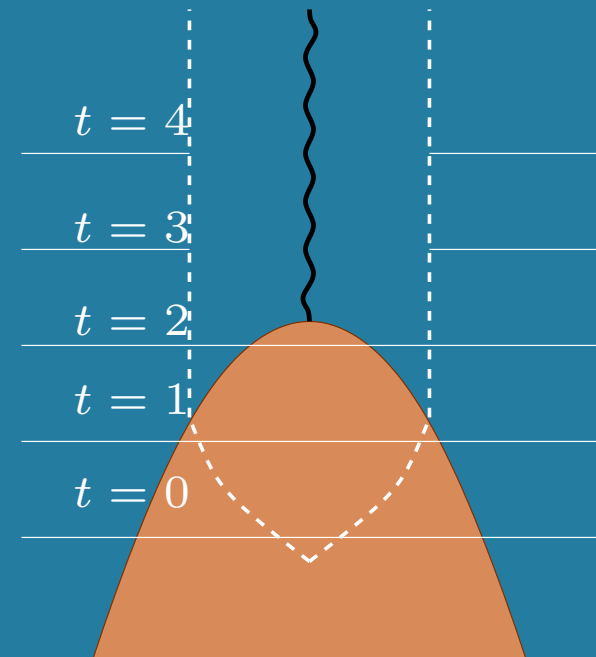
Singularities

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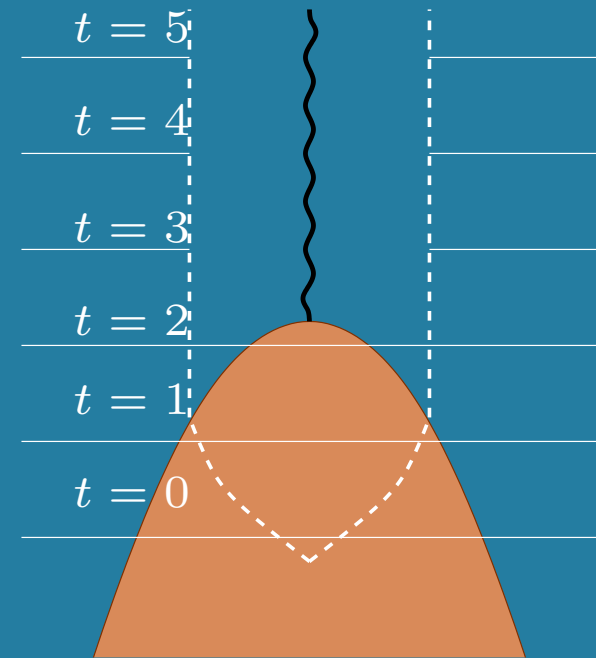
Singularities

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Singularities

- Naive integration hits singularity
- Try coordinate freedom
 - ★ Grid stretching
- Modern method: **Excision**
 - ★ Boundary conditions?
 - ★ Horizon changes shape.

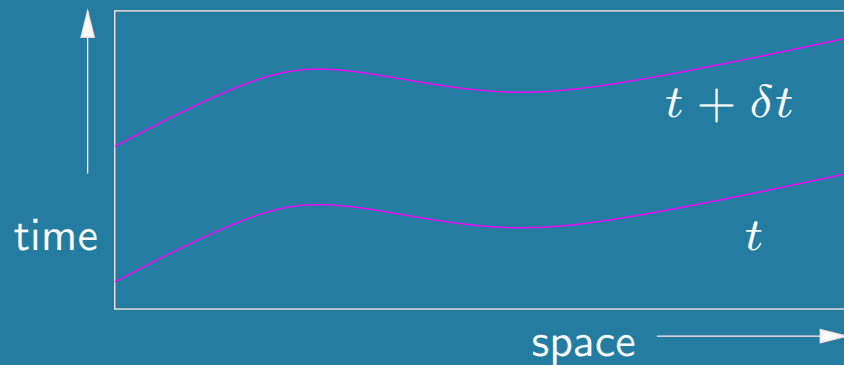


Constraint Violation

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho \quad (C1)$$

Maxwell Illustration: $\partial_t \vec{A} = -\vec{E} - \vec{\nabla} \phi \quad (E1)$

$$\partial_t \vec{E} = \vec{\nabla} \times (\vec{\nabla} \times \vec{A}) - 4\pi\vec{J} \quad (E2)$$



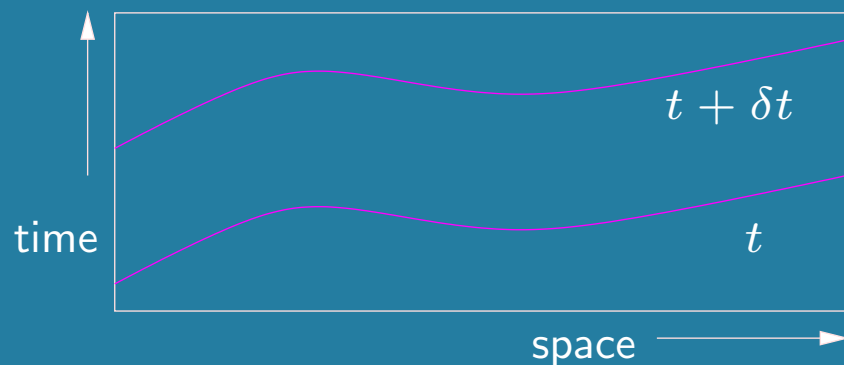
- Does solving (E1), (E2) guarantee (C1)?
 - ★ Exact solution: yes.
 - ★ Approximate solution: no.

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- Does solving (E1), (E2) guarantee (C1)?

- ★ Exact solution: yes.

- ★ Approximate solution: no.

- ★ Problem: 3+1 Einstein evolution equations:

Perturbations that violate constraints can grow rapidly.

Definitions

$$D_i \equiv \partial_i + \text{terms like } g_{jk,i}$$

$$\begin{aligned} \bar{R}_{ij} \equiv & g^{pq} g^{rs} \left(\frac{1}{2} g_{rs,p} g_{q(i,j)} - g_{qs,r} g_{p(i,j)} + \frac{1}{4} g_{ps,j} g_{qr,i} \right. \\ & \left. + g_{pj,s} g_{i[q,r]} + \frac{1}{4} g_{qs,r} g_{ij,p} - \frac{1}{4} g_{rs,q} g_{ij,p} \right) \\ & + g^{pq} \left(g_{p(i,j)q} - \frac{1}{2} g_{ij,pq} - \frac{1}{2} g_{pq,ij} \right) \end{aligned}$$